The authors thank the reviewer for carefully reading our discussion paper and for helpful and constructive comments regarding its content and improvements. We decided to encourage our article with many conclusion both in articles, as in the case of Andreas et al., 2012, Bunker et al., 2003, and during many recent scientific conferences, that further investigation of the differences in the parameterization of the air-sea exchange coefficients are needed.

We would like to start with responding to the major comments. We contain those information also under specific comments.

The aim of the manuscript is to evaluate how much the average monthly and annually momentum transfer values depend on the choice of $C_D$ parameterizations, in other words how the selected parameterization affects the total value of momentum fluxes for large reservoirs. This allows constraining the uncertainty caused by the parameterization choice. In order to achieve this, we used observed wind field for the regions of interest, namely the North Atlantic and the European Arctic, areas where European and Americans oceanographers, including us, operate. This is where most of studies that were basis of the parameterizations we use were performed. We did some comparisons to subtropical basins to see the difference in uncertainty caused by the formula choice between the main study regions and less studied subtopics. In our conclusion, we do not indicate which formula should be used in the future (impossible without new data) in the NA and the EU, but the simple fact that none of the parameterizations used now is final. We don’t want to suggest end users any conclusions because the differences in the parameterizations used are small, and our goal was to help them make an intelligent and deliberate decision about which parameterizations to use.

The text of the review is reproduced below in black type; our comments are in blue; and the changes in the original discussion paper are presented in italics. We reorganized the Introduction as the reviewer suggetg, also removed equation no 7 and reorganized paragraph with this equation, changed the original title to new one and clearly state the purpose of the study.

Evaluation: While each new research on $C_D$ and $\tau$ tries to add a bit more understating and propose improved $C_D(U10)$ expression(s), it is hard to add something truly new after decades of investigations. Thus, it becomes important to clearly state new developments and/or new insights when publishing on this topic. In this sense, this study does not add new knowledge. It is routine. Still, I
see some usefulness of this manuscript in the tabulated $\tau$ differences for different $C_D(U_{10})$ as this can serve as a reference to readers regarding which of many available $C_D(U_{10})$ parameterizations to use. However, the manuscript needs much more work to be suitable for publication. In its present form, it lacks clear objective; could be better organized; and has weak conclusions. I recommend major revisions. Comments and suggestions follow.

We do our best to improve the manuscript, organize it better than it was, clearly state the objective and conclusion, also mark what new it adds to our knowledge. We conclude all of this in our respond and inside the text.

Major comments:

1) The purpose/objective of the manuscript is not clearly formulated. The title suggests two things: a climatology and a review. Two more objectives are hinted in the text (details follow). Focusing on each of these possible objectives would require quite different analyses. With the purpose not well defined, none of the possible objectives is fully developed in the manuscript. Here are specifics.

a) Is the objective a review of $CD$ parameterizations, as the title suggests? If yes, then it is incomplete and without deep physical discussion of the progress and problems regarding parameterizing $CD$. The authors state in Line 43 that it is a must to “take into account other physical processes,” yet they then focus on “wind speed parameterizations, because wind speed is a parameter that is available in every atmospheric circulation model” (Lines 68-69). The authors initiate a review of formulas by dividing them in two groups (Lines 50-55), yet, again, stop short of further discussion on different formulation of roughness length $z_0$. I believe that review and physical discussion on $CD$ are not the motivation of this work. After all, Andreas et al. (2012) and Edson et al. (2013, JPO, DOI: 10.1175/JPO-D-12-0173.1) provide comprehensive recent reviews of the status of parameterizing $CD$.

Yes, you are right. Our aim was not review of Cd parameterizations but checking the uncertainty caused by actual choices done by model and paper authors. That’s why we used only 7 parameterizations, commonly used in literature, or in coupled circulation models. That is why we didn’t do any deep physical discussion for them as the point was to study the spread between the momentum flux values for large basins with realistic wind fields. We did not try to improve the parameterizations using in-situ measurements as this was exactly what Andreas et al. did they calculations using the existing data, and we would simply reproduce their study as we did not have much additional data points. However, we could do something no one has done before, namely check the results of applications
the parameterizations to realistic large-scale wind fields. In fact, we learned that no study did even calculate the average wind stress (momentum transfer) values for most ocean basins since most of the parameterizations were created. This paper meant also to fill that gap, at least partly.

Formulas have been divided into two groups to better familiarize the reader with the used parameterizations. It seems to us that their later detailed description (L73-105) is indirectly explaining the division of the formulas used into two groups.

b) Is the objective a climatology of NA and EA, as the title suggests? If yes, this motivation is not well justified and there is no analysis of the results in climatological terms. If climatology is the objective, then the authors need to tell us why they focus on the NA and EA regions? What atmospheric and oceanic conditions does the CD parameterization need to represent well in these regions? If climatology is the goal, what is the temporal or spatial reference? It seems the chosen spatial references are the global ocean and the Tropics, to which the authors compare their results for NA and EA. But if $\tau$ is obtained globally and over many regions (Table 1), then why emphasize NA and EA in the title? Why look into differences due to CD formulation between northern regions and the Tropics, when it is certainly expected to have differences due to geography? As for a temporal reference, the authors should choose a period within or outside the 1992-2010 period which gives average atmospheric and oceanic conditions, not affected by long-term variations such as the North Atlantic Oscillation, which changes the position of jet stream and thus the wind and SST fields at the surface; these, in turn, change the wind stress. If climatology is the goal, then the authors should analyze their $\tau$ results for trends and variations over the 1992-2010 period. Should give annual as well as inter-annual variations. Finally, in my opinion, to provide a comprehensive regional climatology, the authors should analyze long-term $\tau$ values obtained with one, chosen CD formulation in order to clearly isolate climatologically-relevant variations.

Our aim was not a climatology of the NA and the EA. The original title was wrongly defined and we change it. The main areas of the study: North Atlantic and European Arctic seas are the areas where European and Americans oceanographers, including us, operate. This is where most of studies that were basis of the parameterizations we use were performed. We did some comparisons to sub-tropical basins to see the difference in uncertainty caused by the formula choice between the main study regions and less studied subtopics.

c) Is the objective to evaluate CD parameterizations and recommend a new one for use in circulation coupled models (Line 136)? A hint for such an objective comes from the authors’ conclusion "the parameterizations used
in the models possibly need upgrading” (Lines 332-333). If this is the objective, then the authors should give us a list of parameterizations used in different circulation models; discuss the advantages and limitations of these currently-used $CD$ parameterizations; then demonstrate how other $CD$ parameterizations would do better. For climate and circulation models, the $CD$ parametrization is important for the mixing layer depth. So the authors should show how new $CD$ parameterization would improve the modeling of the mixing layer. The manuscript offers limited information on what the current models use (Lines 136-140). There is no analysis on how $CD$ would affect the performance of model variables related to $CD$. So the conclusion in Lines 332-333 is not convincing for modelers.

It wasn’t our objective at this study/manuscript. We used the sentence “the parameterizations used in the models possibly need upgrading” because we want to show that there is still a lot to do and there is a room for it, despite widespread statements that it is hard to add something truly new after decades of investigations. We decided to encourage our article with many conclusion both in articles, as in the case of Andreas et al., 2012, Bunker et al., 2003, and during many recent scientific conferences, that further investigation of the differences in the parameterization of the air-sea exchange coefficients are needed. However, we are agree that this sentence was not the best so we improved it: L326 used in the models possibly need further development.

d) Is the objective to demonstrate/justify the need of new measurements in NA and EA for improved $CD$ parameterization in high latitudes? A hint for such an objective comes from Lines 188-189 regarding frequent ship deployment in EA, including “R/V Oceania, the ship of the institution the authors are affiliated with.” If this is the objective, the manuscript would take completely different direction with discussion and analysis related to measuring methods and quality of data necessary for $CD$. Of course, this is not the objective because the authors say in Lines 65-68 that their intention “is not to re-invent or formulate a new drag parameterization … but to revisit definition of the existing drag parameterization.”

The aim of the manuscript is to evaluate how much the average monthly and annually momentum transfer values depend on the choice of $C_D$ parameterizations, in other words how the selected parameterization affects the total value of momentum fluxes for large reservoirs. This allows constraining the uncertainty caused by the parameterization choice. In order to achieve this, we used observed wind field for the regions of interest, namely the North Atlantic and the European Arctic.

As the reviewer pointed out that the aim of the manuscript was not clearly formulated, so we changed it and added properly information and corrections
Confirming and explaining the nature and consequence of the interaction between the atmosphere and ocean is one of the great challenges in climate and sea research. These two spheres are coupled which lead to variations covering time scales from minutes to even millennia. The purpose of this article is to examine, using a modern set of software processing tools called the FluxEngine, the nature of the fluxes of momentum across the sea surface over the North Atlantic and the European Arctic. These fluxes are important to determine current system and sea state conditions. Our goal is to evaluate how much the average monthly and annually momentum transfer values depend on the choice of $C_D$, using the actual wind field from the North Atlantic and the European Arctic, and demonstrate existing differences as a result of the formula used.

In this paper we investigate how the relevant or most commonly used parameterizations for drag coefficient ($C_D$) affect the value of momentum transfer values, especially in the North Atlantic (NA) and the European Arctic (EA). Our task was to demonstrate existing differences as a result of the formula used how big they can be. As is widely known, the exact equation that describes the connection between the drag coefficient and wind speed depends on the author (Geernaert, 1990). Our intention here is not to re-invent or formulate a new drag parameterization for the NA or the EA, but to revisit the existing definition of drag parameterization, and, using satellite data, to investigate how existing formulas represent the environment in the North. We concentrated on wind speed parameterizations, because wind speed is a parameter that is available in every atmospheric circulation model. Therefore, it is used in all air-sea flux parameterizations, and presently it is used even when sea state provides a closer physical coupling to the drag coefficient (for review see Geernaert et al., 1986).

The original title could cause confusion and did not clearly define the paper purpose. Therefore we change it in the revised version. The new title now is:

**Effect of drag coefficient formula choice on wind stress climatology in the North Atlantic and the European Arctic**

e) I am listing all these possible objectives only to make the point that the authors need a well stated objective in order to focus their analysis and discussion.
I believe the authors wish to assess $CD$ parameterizations only to decide which
one to use in some larger project.

Yes, you are indirectly right.

For such an assessment, the authors only need to clearly tell us why they consider the formulations (8)-(14) (i.e., no need of comprehensive review). They do not need climatology to make this assessment. One year data for \( t \) is enough.

We have chosen the formulations 7-13 as, in our opinion, they are most common used ones in the literature. The results between the formulae used can came from different functional \( C_D \) formulation as well as seasonal variations in NA and EA. In conclusion we add information about it.

However, to make the decision, the authors need not only to quantify differences (this is what the current results offer). They need also to make a thorough analysis what causes the differences. They need to evaluate how much of the differences come from: (i) different functional \( C_D \) formulation; (ii) different quality of the data on which the parametrizations are based; and (iii) seasonal variations in NA and EA.

Differences in momentum flux mostyl came from the different functional \( C_D \) formulation than from seasonal variations in the NA and the EA.

The authors also need some reference to show them which \( C_D \) formulation is suitable for NA and EA. Perhaps comparison of their results to regional data? Perhaps an investigation of how well a feature specific to NA and EA is represented when using different \( C_D \) parameterizations? With such direction of the manuscript, the title may need revision to exclude claims on climatology and review.

We believe that there is no answear for that in the literature. If there was one, our paper would be pointless. We can only guess that the newer parameterizations, based on more observations, are better but this is only a guessing. Our paper tries to answer a different question: what are the differences between those parameterizations when applied to obseved wind fields in a given basin. This is the question of uncertainty due to the choice of parameterization. We also do not agree that an analysis of which parameterization has lower or higher values for which winds would relly help. It is a trivial observation that those with highest values for strong winds give the highest wind stressess (with the exception of Andreas et al., which partly offsets that with low values for low winds). We added the following sentence at line 258.
In the NA region with winds stronger than average for world ocean, the formula giving highest momentum transfer results are the ones with highest values for strong winds, with exception of Andreas et al. (2012) which is lower due to its low values for lower winds speeds.

2) There are several typos in the formulae that need fixing. Most importantly, it is necessary to check the coding for the calculations. These typos are as follows. In (2), \( U_{102} \) is in the denominator. In (7), why \( U_{10} \) is squared? The relationship between \( u^* \) and \( U_{10} \) is linear (Andresa et al., 2012, their eq. 1.10; Edson, et al., 2013, their eq. 22). In (14), need square on the wind ratio \( (u^*/U_{10N})^2 \); in the parenthesis, \( U_{10N}^2 \) is in the denominator; needs square on the parenthesis (compare to Andreas et al., 2012, their eq. 1.10).

We apologize for all errors in the manuscript. We checked all formulas again and corrected the mistakes.

3) Give better justification on choosing \( CD \) parameterizations (8) \( \& \) (14). For example, it seems you have chosen \( CD \) parameterizations formulated as power law, linear, polynomial, constant. Why do you need (9) and (10)? They are so similar? Describe the merits of (11), (13) and (14), as well as their differences (e.g., data on which they are based). Do these formulations account implicitly for different processes in addition to \( U_{10} \)?

All parameterizations are important to the history of the field and all have been widely used by other authors. In our opinion, skipping any of them, especially the oldest and newest would cause the study flawed. In fact, if we had to change the number of formulas, we would rather increase it than decrease. This number was a compromise to make presenting the results graphically not too overwhelming for the reader.

We have chosen parameterization no 9 and 10 because despite the fact that the formulas themselves are so similar and have the same source (based on Charnock’s relations) Garratt in his research showed that this formula is suitable only for lights wind (over 4 m s\(^{-1}\)), while Wu based on this statement showed that this formula can be proposed for all sea state and fit closely to the data throughout the entire wind-velocity range. Both formulas were used in the literature.

Yelland and Taylor used an automatic inertial dissipation system, over the Southern Ocean, to obtain data for wind stress estimations. During their study they examined the balance between local production and dissipation of turbulent kinetic energy. It is the newest version of the linear parameterizations and we wanted to check (and show) how much difference in integrated momentum flux the three of them make in comparison to the other
parameterizations.

Large and Yeager parameterizations came from a compilation of global data sets from different sources, like NCEP/NCAR, CCSM, historical SST, and it was developed for configurations OGCM and coupled OGCM-SIM models. They used observation for winds from 1 to more than 25 m/s speed. It is used in many modern coupled circulation models.

Andreas et al., used data from the literature to test approach proposed by Foreman and Emeis based on the eddy-covariance flux measurement over the sea to deduce air-sea drag relations. For their study their used data for very strong winds (>24 m/s). We believe this is the most up to date parameterization and on the other hand not well known and appreciated.

All of them used neutral-stability wind speed

L101-108 All of them are generated from the vertical wind profile, but they differ in the formulas used. Two of the parameterization which we chosen are formulated as power-law of the relationship between $C_D$ and $U_{10}$ (eq. 7 and 13), three are formulated as linear-law (eq. 8, 9 10 for light winds, and 12), and one as constant value of the relationship (eq. 11). All the above studies propose different parameterizations (see Fig. 1) of the drag coefficient and the function of wind speed, which reflects the difficulties in simultaneously measuring at high sea stress (or friction velocity) and wind speed.

4) Suggest re-organizing the Introduction to include Lines 37-72 plus one paragraph on why you focus on NA and EA, then another paragraph clearly stating the objective of the study. Suggest combining Lines 73-154 with Lines 195-215 in one section dedicated on $CD$ parameterizations. Only parts of the historical (incomplete) review in lines 73-154 are necessary. Start with the definitions in Lines 73-93. Then introduce (8) \( \text{(14)} \) one by one. Add information on MOST (Lines 115-122) and circulation models (Lines 136-140) only when they are needed, e.g., when you introduce (11) and (12), respectively. Finish the section with Lines 155-160. With this organization you will avoid the current inconsistency of presenting Fig. 1 with all parameterizations before they are described. Remove lines 122-127 and Lines 130-135 because you do not use Trenberth et al. (1989) and COARE algorithm. Unless you decide to use COARE 3.5 as a reference.

Done

L23-145 Confirming and explaining the nature and consequence of the interaction between the atmosphere and ocean is one of the great challenge in climate and sea research. These two sphere are coupled which lead to variations covering time scales from minutes to even millennia. The purpose
of this article is to examine, using a modern set of software processing tools called the FluxEngine, the nature of the fluxes of momentum across the sea surface over the North Atlantic and the European Arctic. These fluxes are important to determine of current system and sea state conditions. Our goal is revisit how the existing definition of drag parameterization affects the value of total momentum fluxes, using the actual wind field from the North Atlantic and the European Arctic and demonstrate existing differences as a result of the formula used.

The ocean surface mixed layer is a region where kinematic forcing affects the exchange of horizontal momentum and controls transport from the surface to depths (Gerbi et al., 2008, Bigdeli et al., 2017). Any attempt to properly model the momentum flux from one fluid to another as the drag force per unit area at the sea surface (surface shear stress, $\tau$) take into account other physical processes responsible for generating turbulence such as boundary stress, boundary buoyancy flux, and wave breaking (Rieder et al., 1994, Jones and Toba, 2001). Fluxes across the sea surface usually depend nonlinearly on the relevant atmospheric or oceanic parameters. Over the past fifty years, as the collection of flux data has increased, many empirical formulas have been developed to express the ocean surface momentum flux as a relationship between non-dimensional drag coefficient ($C_D$), wind speed ($U_{10}$), and surface roughness ($z_0$) (Wu 1969, 1982; Bunker, 1976; Garratt, 1977; Large and Pond, 1981; Trenberth et al., 1989; Yelland and Taylor, 1996, Donelan et al., 1997; Kukulka et al., 2007; Andreas et al., 2012). These formulas can be divided into two groups. One group of theories gives the $C_D$ at level $z$ in terms of wind speed and possibly one or more sea-state parameters (for example, Geernaert et al., 1987, Yelland and Taylor, 1996, Enriquez and Friehe, 1997), while the second group provides formulas for roughness length $z_0$ in terms of atmospheric and sea-state parameters (for example, Wu, 1969, Donelan et al., 1997, Andreas et al., 2012 (further referred to as A12)).

As the exchange of air-sea momentum is difficult to measure directly over the ocean meteorologist and oceanographers often rely on bulk formulas parameterized by Taylor (1916), that relate the fluxes to averaged wind speed through transfer coefficients:

$$\tau = \rho C_{Dz} U_z^2$$  \hspace{1cm} (1)
where $\tau$ is the momentum flux of surface stress, $\rho$ is air density, $C_{Dz}$ is the non-dimensional drag coefficient appropriate for $z$ height, and $U_z$ is the average wind speed at some reference height $z$ above the sea. $C_{Dz}$ is commonly parameterized as a function of mean wind speed ($m \, s^{-1}$) for neutral-stability at a 10 m reference height above mean sea level (Jones and Toba, 2001), which is identified as $C_{DN10}$ or $C_{D10}$ (this permits avoiding deviation for the vertical flow from the logarithmic law):

$$C_{DN10} = \frac{\tau}{\rho U_{10}^2} = \left(\frac{u^*}{U_{10}}\right)^2$$

(2)

where $u^*$ is friction velocity. Alternatively, the neutrally stratified momentum flux can be determined from the logarithmic profile, thus Eq. 1 can be express as:

$$C_{DN10} = \left[\kappa \ln (10/z_0)\right]^2$$

(3)

where $z_0 (m)$ is the aerodynamic roughness length, which is the height, above the surface to define the measure of drag at which wind speed extrapolates to 0 on the logarithmic wind profile (Andreas et al., 2012), and $\kappa$ is von Kármán constant ($\kappa=0.4$).

At the same time, we can define the value of friction velocity by the following equation:

$$\tau = \rho u^2$$

(4)

Comparison with bulk formula (1) leads to the equation:

$$u_z^2 = C_{D10} U_{10}^2$$

(5)

Some of the first studies (Wu, 1969, 1982, Garrat, 1977) focused on the relationship between wind stress and sea surface roughness, as proposed by Charnock (1955), and they formulated (for winds below 15 m $s^{-1}$) the logarithmic dependence of the stress coefficient on wind velocity (measured at a certain height) and the von Kármán constant. Currently common parameterizations of the drag coefficient are a linear function of 10 m wind speed ($U_{10}$), and the parameters in the equation are determined empirically by fitting observational data to a curve. The general form is expressed as (Guan and Xie, 2004):
\[ C_D10^3 = (a + bU_{10}) \]  

In this work our focus is on the fluxes of average values using seven different drag coefficient parameterizations \(C_D\), chosen for their importance for the history of the field out of many published within the last half century (Bryant and Akbar, 2016).

\[ 10^3 \cdot C_{D10} = 0.5U_{10}^{0.5} \quad \text{for} \quad 1 \text{ m s}^{-1} < U_{10} < 15 \text{ m s}^{-1} \]  

(Wu, 1969)

\[ 10^3 \cdot C_{DN10} = 0.75 + 0.067U_{10} \quad \text{for} \quad 4 \text{ m s}^{-1} < U < 21 \text{ m s}^{-1} \]  

(Garratt, 1977)

\[ 10^3 \cdot C_{D10} = (0.8 + 0.065U_{10}) \quad \text{for} \quad U_{10} > 1 \text{ m s}^{-1} \]  

(Wu, 1982)

\[ 10^3 \cdot C_{DN10} = 0.29 + \frac{3.1}{U_{10N}} + \frac{7.7}{U_{10N}^2} \quad \text{for} \quad 3 \text{ m s}^{-1} < U_{10N} < 6 \text{ m s}^{-1} \]  

(Yelland and Taylor, 1996)

\[ 10^3 \cdot C_D = 1.3 \quad \text{everywhere} \]  

(NCEP/NCAR)

\[ 10^3 \cdot C_{DN10} = \frac{2.7}{U_{10N}} + 0.142 + 0.076U_{10N} \quad \text{everywhere} \]  

(Large and Yeager, 2004)

\[ C_{DN10} = \left(\frac{u^*}{U_{10N}}\right)^2 = a^2 \left(1 + \frac{b}{a} U_{10N}\right)^2 \quad \text{everywhere} \]  

(Andreas et al., 2012)

where \(C_{DN10}\) is the expression of neutral-stability (10-m drag coefficient), \(C_{D10}\) is the drag coefficient dependent on surface roughness, \(U_{10}\)
is the mean wind speed measured at 10 m above the mean sea surface, $U_{10N}$ is the 10-m, neutral-stability wind speed. All of them are generated from the vertical wind profile, but they differ in the formulas used. Two of the parameterization which we chosen are formulated as power-law of the relationship between $C_D$ and $U_{10}$ (eq. 7 and 13), three are formulated as linear-law (eq. 8, 9 10 for light winds, and 12), and one as constant value of the relationship (eq. 11). All the above studies propose different parameterizations (see Fig. 1) of the drag coefficient and the function of wind speed, which reflects the difficulties in simultaneously measuring at high sea stress (or friction velocity) and wind speed.

Wu (1969), based on data compiled from 12 laboratory studies and 30 oceanic observations, formulated power-law (for breezes and light winds) and linear-law (for strong winds) relationships between the wind-stress coefficient ($C_y$) and wind velocity ($U_{10}$) at a certain height $y$ at various sea states. In his study, he used roughness Reynolds numbers to characterize the boundary layer flow conditions, and he assumed that the sea surface is aerodynamically smooth in the range of $U_{10} < 3 \text{ m s}^{-1}$, transient at wind speed $3 \text{ m s}^{-1} < U_{10} < 7 \text{ m s}^{-1}$, and aerodynamically rough at strong winds $U_{10} > 7 \text{ m s}^{-1}$. He also showed that the wind-stress coefficient and surface roughness increase with wind speed at light winds ($U_{10} < 15 \text{ m s}^{-1}$) and is constant at high winds ($U_{10} > 15 \text{ m s}^{-1}$) with aerodynamically rough flow. Garratt (1977), who assessed the 10 m neutral drag coefficient ($C_{DN10}$) based on 17 publications, confirmed the previous relationship and simultaneously suggested a linear form of this relationship for light wind. Wu (1980) proposed the linear-law formula for all wind velocities and later (Wu, 1982) extended this even to hurricane wind speeds. Yelland and Taylor (1996) presented results obtained from three cruises using the inertial dissipation method in the Southern Ocean and indicate that using the linear-law relationship between the drag coefficient and wind speed (for $U_{10} > 6 \text{ m s}^{-1}$) is better than using $u^*$ with $U_{10}$. The NCEP/ NCAR reanalysis (Kalnay et al., 1996) uses a constant drag coefficient of $1.3 \times 10^{-3}$ while, for example, the Community Climate System Model version 3 (Collins et al., 2006) uses a single mathematical formula proposed by Large and Yeager (2004) for all wind speeds. Andreas et al. (2012) based on available datasets, friction velocity coefficient versus neutral-stability wind speed at 10 m, and sea surface roughness tested the approach proposed by Foreman and Emeis (2010) for friction velocity in order to find the best fit for parameters $a =$
0.0583 and \( b = -0.243 \). They justify their choice by demonstrating that \( u^* \) vs. \( U_{10N} \) has smaller experimental uncertainty than \( C_{DN10} \), and that one expression of \( C_{DN10} \) for all wind speeds overstates and overestimates results in low and high winds (Figs. 7 and 8 in A12).

In this paper we investigate how the relevant or most commonly used parameterizations for drag coefficient (\( C_D \)) affect to value of momentum transfer values, especially in the North Atlantic (NA) and the European Arctic (EA). Our task was to demonstrate existing differences as a result of the formula used how big they can be. As is widely known, the exact equation that describes the connection between the drag coefficient and wind speed depends on the author (Geernaert, 1990). Our intention here is not to reinvent or formulate a new drag parameterization for the NA or the EA, but to revisit the existing definition of drag parameterization, and, using satellite data, to investigate how existing formulas represent the environment in the North. We concentrated on wind speed parameterizations, because wind speed is a parameter that is available in every atmospheric circulation model. Therefore, it is used in all air-sea flux parameterizations, and presently it is used even when sea state provides a closer physical coupling to the drag coefficient (for review see Geernaert et al., 1986).

5) Section 3 “Result” is straightforward. It describes Table 1, maps, and seasonal graphs. To make these results useful, you need to extend the analysis of these data, discuss what causes the differences; and suggest which \( C_D \) parameterizations is useful for NA and EA.

It is impossible to tell which formula is better comparing the results of its integration with wind field. That’s is why we study mainly the spread (uncertainty) of result coming from the parameterization choice. However the most recent one (Andreas et al., 2012) is based on the largest measurement set, and that’s why we point out what it may imply for momentum transfer, especially at low wind speed. We adds properly information inside the text.

We add properly new information inside the text at Result section:

L190-198 Despite many measurements, the drag coefficient still has wide variability at low and moderate wind speeds. Our research has showed that all lower wind values (<10 m s\(^{-1}\)) the differences between the drag coefficient parameterizations are greater than at higher speeds (> 10 m s\(^{-1}\)) and the most outlier results are those obtained from the power law parameterization
of Andreas et al., (2012), which are characterized by a sinusoidal distribution relative to the wind speed. The lower the wind speed, the higher uncertainty are, and at low winds it is uncertainty by a factor of 0.5-1.5 depending on the formula used, while at moderate winds it is uncertainty by a factor of 1.5-2.0 (Fig. 1).

We improved also the conclusion by adding some new information.

L288-326 In the present work the evaluation of how the selected parameterization affects the total value of momentum fluxes for large reservoirs was assessed. This allows constraining the uncertainty caused by the parameterization choice. In order to achieve this we calculated monthly and annually average momentum fluxes using a set of software processing tools called the FluxEngine in the North Atlantic (NA) and the European Artic (EA). The NA was defined as all sea surface areas in the Atlantic sector north of 30°N, and the EA was sea areas north of 64°N. Based on our results, we still do not know which one of the parameterizations can be recommend as the most suitable for the NA and the EU study. Further investigation of the differences in the parameterization of the exchange coefficient in the various algorithm would help in resolving this problem. Despite many measurements, the drag coefficient still has wide variability at low and moderate wind speeds. The lower the wind speed, the higher uncertainty are, and at low winds it is uncertainty by a factor of 0.5-1.5 depending on the formula used, while at moderate winds it is uncertainty by a factor of 1.5-2.0 (Fig. 1). The annual mean wind speed in the NA is 10 m s⁻¹, and in the EU it is 8.5 m s⁻¹.

We show that the choice of drag coefficient parameterization can lead to significant differences in resultant momentum flux (or wind stress) values. Comparing the values of momentum flux across the sea surface from the power law parameterization, it showed that in both regions, with low and high winds, the parameterizations specified for all winds speeds (eq. 13) has lower values of wind stress than the parameterizations specified for light winds (eq. 7). In the Arctic, the NA, and globally the differences between the wind stress, depend on formula used, are 14 % and they are higher in low winds areas. The parameterizations generally have a decreasing trend in the resultant momentum flux values, with the most recent (Andreas et al., 2012) producing the lowest wind stress values, especially at low winds, resulting in almost 20 % differences in the tropics (Table 1). The differences can be much larger on monthly scales, up to 29 % in the NA and 36 % in the EA (in
months of low winds) and even 50% locally in the area west of Spitsbergen. For months that have the highest average winds, the percentage differences are smaller (about 7% everywhere), but because absolute value of the flux are largest for high winds, this 7% discrepancy is also important for air-sea momentum flux values. Since momentum flux is an important parameter in ocean circulation modeling, we believe more research is needed (one aspects that needs more research is the fact that the newest power law parameterization, A12, produces less momentum flux than all the previous ones, especially in lower winds), and the parameterizations used in the models possibly need further development.

Additional comments:

Title: If possible (perhaps talk with the OS editor), revise the title to better reflect the purpose of your manuscript.
We changed the title. Now it is:

Effect of drag coefficient formula choice on wind stress climatology in the North Atlantic and the European Arctic

Abstract: Too long, dilutes what you did and what you have found. Suggest substantial shortening. Avoid giving references in the abstract. Refer to different CD parameterizations by their specific characteristics (e.g., power law, linear, etc), not by author.
Done.

Lines 18-19 and Lines 227-230: Oldest vs newest CD parameterization. This is not the most important difference. Frame your discussion around the functional form, the data they are based on, how well they represent low and high wind conditions.

Line 14-18: When we choose the parameterizations which increased linearly with wind speed (7-9) momentum flux were largest for all months, in compare to values from the two parameterizations which increase with wind speed sinusoidal (12 and 13) in both regions with high and low winds and CD values were consistently higher for all wind speeds.

Line 186-190: The differences between the parameterizations are distinct (Fig. 1). The CD values from the parameterizations 7-9 increased linearly
with wind speed since the results from the parameterizations 10,12,13 are characterized by sinusoidal distribution and indicating decreases for winds in the range of 0 - 10 m s\(^{-1}\), after which they began increase.

Line 22-23: Suggest removing this sentence. This is common sense, no need to be in the abstract.

We removed all the sentences from lines 22 to 27: For global data not much seasonal change was note due to the fact that the strongest winds are in autumn and winter as these seasons are inverse by six months for the northern and southern hemispheres. The situation was more complicated when we considered results from the North Atlantic, as the seasonal variation in wind speed is clearly marked out there. With high winter winds, the A12 parameterization was no longer the one that produces the smallest wind stress.

Line 30: “the sequence of values” is the least important thing to discuss about the differences. Discuss the physical behavior.

We removed sentences at lines 28-30: However, for low summer winds, it is the lowermost outlier. As the A12 parameterization behaves so distinctly differently with low winds, we showed seasonal results for the tropical ocean. The sequence of values for the parameterization was similar to that of the global ocean, but with visible differences between NCEP/NCAR, A12 and LY04 parameterizations. Because parameterization is supported with the largest experimental data set observations of very low (or even negative) momentum flux values for developed swell and low winds, our results suggest that most circulation models overestimate momentum flux.

and reorganized the rest:

L18-21 As the one of power law parameterization (13) behaves so distinctly differently with low winds, we showed seasonal results for the tropical ocean, which were subdued for the whole region, with monthly averages in the range of 0.2 to 0.3 N m\(^2\).

Line 76: Definition of \(\tau\) is already given in Lines 42-43. Here, and many other places, remove repeated definitions. Done.

Lines 89-90: Suggest removing this sentence, repeats definition given in Line 83.

We reorganized this sentence:
At the same time, we can define the value of friction velocity by the following equation:

Line 140: I guess you mean here equation (5), which assumes proportionality; (6) modifies (5) to linear relationship.
Line 144: I think you mean here equation (7), not (8). Eq. (7) needs correction (see Major comment 2).
Line 147: Fix symbol UN10 to $U_{10N}$. Check all your math symbols for correctness and consistency.

We reorganized this sentence and remove some of the information from them because it seemed unnecessary after the reorganization:

Andreas et al. (2012) based on available datasets, friction velocity coefficient versus neutral-stability wind speed at 10 m, and sea surface roughness tested the approach proposed by Foreman and Emeis (2010) for friction velocity in order to find the best fit for parameters $a = 0.0583$ and $b = -0.243$. They justify their choice by demonstrating that $u^*$ vs. $U_{10N}$ has smaller experimental uncertainty than $C_{DN10}$ and that one expression of $C_{DN10}$ for all wind speeds overstates and overestimates results in low and high winds (Figs. 7 and 8 in A12).

Line 155: Fig 1 shows parameterizations whose equations are not yet introduced. Need to introduce (8) $f(14)$ before referring to Fig. 1. See Major comment 4.

We divided these sentences. The part was moved to lines 105-108 and part to line 133-135.

All the above studies propose different parameterizations (see Fig. 1) of the drag coefficient and the function of wind speed, which reflects the difficulties in simultaneously measuring at high sea stress (or friction velocity) and wind speed.

In this paper we investigate how the relevant or most commonly used parameterizations for drag coefficient ($C_D$) affect to value of momentum transfer values, especially in the North Atlantic (NA) and the European Arctic (EA).

Line 168: Use symbol $U_{10}$ instead of re-defining it again.
Done

Line 168-169: How these data on sea roughness are used? None of your equations (8) $f(14)$ uses sea roughness. Why then introduce these data here?
Our mistake. Unnecessarily and wrongly introduced the sea surface roughness.
Line 175: Use symbol $U10N$ instead of re-defining it again.
Done

Lines 177-179: You do not use wave data in (8) $\xi(14)$, why do you introduce these data here?
Also our mistake. We have also introduced it unnecessarily.

Lines 180-183: Are all these details part of the FluxEngine software? Or are these done by you?
All of these details are already part of the FluxEngine. We have added relevant information to the text and reference to the literature:

L161-164 The data layers within each output file, which are details part of the FluxEngine, include statistics of the input datasets (e.g., variance of wind speed, percentage of ice cover), while the process indicator layers include fixed masks as land, open ocean, coastal classification, and ice.

Lines 186, 226, 235: Suggest re-numbering Fig. 6 to Fig. 2, then all other figures. You refer to all other figures much later in the text.
Done

Lines 195-215: Need to be introduced before Line 155 (see Major comment 4).
Done

Line 217: “gridded global air-sea momentum” Why global when your emphasis is on NA and EA? Is global a good reference? You need representation of average conditions (either spatially or temporally averaged) for a reference. Need to work this out.

FluxEngine software produced only global fluxes grid data and after that we calculated monthly values for separated region.

L177-180 Using the FluxEngine software, we produced global gridded monthly air-sea momentum fluxes and from these we have extracted the values for the study region, the global ocean, the NA Ocean, and its subsets: the Arctic sector of the NA and the West Spitsbergen area (WS).

Line 229: Revise “sinusoidal”. The decrease at low winds is not due to sinusoidal behavior.

Line 187-190: The $C_D$ values from the parameterizations 7 - 9 increased linearly with wind speed since the results from the parameterizations
10,12,13 are characterized by sinusoidal distribution and indicating decreases for winds in the range of 0 - 10 m s$^{-1}$, after which they began increase.

Lines 246-248: Why looking into global values for seasonal variations when it is clear that opposite seasons cancel the variations? For seasonal variations, it is better to compare to Northern (or Southern) hemisphere.

We done this as we want to showed results from regionally scale against the larger background and to show the order of magnitude of differences in Northern hemisphere, and also for better detail results from regionally scale.

Line 276: “could be at statistical effect” What do you mean? Suggest revision for clarity.

We revised that to „an averaged effect”. We meant that sub-tropical trade wind areas tend to have stable winds of speeds for which the Andreas et al. (2012) parameterization has almost no drag which is due to waves and wind travelling at similar velocities. We reorganized paragraph with this sentence:

L201-208 For this estimation we chose the two most-recent parameterizations (eq. 12 and 13) that showed the lowest values and change seasonally depending on the area used. As a result, these months with weak winds have significantly lower momentum flux values, which could be the effect of statistically weaker wind in ocean areas having stable winds with waves traveling in the same direction as the wind at similar speeds. Comparison showed that the A12 parameterization demonstrates almost zero sea surface drag for winds in the range of 3 - 5 m s$^{-1}$, which is compensated for by a certain surplus value for strong winds.


The annual ratios of the parameterizations. We change this in the text.

L254-256 A surprising result is the annual ratios of the parameterizations values for the global, the NA, and the Arctic regions (Fig. 4 shows that this is not true on monthly scales).

Lines 329-331: Not clear what is your conclusion here. Please revise.

The sentence is now:

L319-322 For months that have the highest average winds, the percentage differences are smaller (about 7 % everywhere), but because absolute value of the flux are largest for high winds, this 7% discrepancy is also important
Line 333: “need upgrading” From what expression? To what expression? You make all these calculations but in the end you do not recommend what is good to use. See Major comment 1c.

Changed to „need further improvements“.

We do not indicate which formula should be used in the future (again impossible without new data) in the NA and the EU, but the simple fact that none of the parameterizations used now is final. We don’t want End Users to draw any conclusions because the differences in the parameterizations used are small, and our goal was to help them mak an intelligent and deliberate decision about which parameterizations to use.

L322-326 Since momentum flux is an important parameter in ocean circulation modeling, we believe more research is needed (one aspects that needs more research is the fact that the newest power law parameterization, A12, produces less momentum flux than all the previous ones, especially in lower winds), and the parameterizations used in the models possibly need further development.


Sorry for that, of course area average. What we had in mind in this table and fig. 5 was the average annual value of the momentum flux divided by the surface value of each area.

We have changed it and revised in the text.

L488-490 and 514-516 Table 1. Area average annual mean values of momentum flux (wind stress) [N m⁻²] for all the studied regions and parameterizations. In each column the percentage values are normalized to A12, the parameterization that produced the smallest average flux values.

L506-510 and 657-661 Figure 5. Area annual average momentum flux values for (a) European Arctic and (b) Tropical ocean. The vertical solid line is the average of all seven parameterization and the dashed lines are standard deviations for the presented values. Global and the North Atlantic results are not shown because the relative values for different parameterizations are very similar (see Table 1), scaling almost identically between the basins.

L248-256 Table 1 and Fig. 5 present the annual average air-sea momentum flux values (in N m⁻²) for all the all regions studied and all the
parameterizations. The results show that the annual North Atlantic momentum fluxes, depending on the formula used, varies from $-0.0.290 \text{ N m}^{-2}$ for A12 to $0.333 \text{ N m}^{-2}$ for Wu (1969). In the case of global annual average, the values are -0.283 and 0.322, respectively. Table 1 shows also the same data “normalized” to the A12 data (presented as percentages of A12, which produced the lowest values for each region), which allows us to visualize the relative differences. A surprising result is the annual ratios of the parameterizations values for the global, the NA, and the Arctic regions (Fig. 4 shows that this is not true on monthly scales).

Figure 3a: Should show data for the Northern hemisphere if you want to use this as a reference for seasonal variations.
Yes, as in our study we compare globally data with data from Norther hemisphere. They are important as a references.

Fig. 4: Why do you need this figure? What more does it shown than Fig. 1?

We have included this figure to better illustrate the differences between the two important parameterizations. We thought about the reviewer’s comment, which we agree with, therefore we decided to remove this chart as it adds nothing new to the article. The descriptions of fig. 4 from lines 270-282 were reorganized and joined with the descriptions of fig. 1. Line 192-203

Writing style and corrections:

Line 32: I guess you mean “Because A12 parameterization..”
We removed this sentence.
Line 37: Suggest revision to read: “Wind stress at the air-sea interface influences the wind-wave interaction, including…”
Done.
Line 43: Suggest revising “must” to other word. “Must” is a firm request, which you do not follow in your subsequent considerations.
Done.
Lines 45-46: Suggest revising to read: “…fifty years, as the collection of flux data has increased, many empirical formulas…”
Done.
Line 61: Suggest revising “we chose to check” to “we investigate” or “we quantify”
Done.
Line 67: Suggest revising “accommodate” to “represent”
Done.
Lines 82-83: You have $u^*$ in italic and non-italic. Here, and everywhere, give
mathematical symbols consistently.

Done.

Line 144: Abbreviation A12 should be introduced on first encounter, in line 50.

Done.
Effect of drag coefficient formula choice on wind stress climatology in the North Atlantic and the European Arctic

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Key points: drag coefficient, European Arctic, North Atlantic, parameterizations

Abstract

In this paper we have chosen to check the differences between the relevant or most commonly used parameterizations for drag coefficient ($C_D$) for the momentum transfer values, especially in the North Atlantic (NA) and the European Arctic (EA). We studied monthly values of air-sea momentum flux resulting from the choice of different drag coefficient parameterizations, adapted them to momentum flux (wind stress) calculations using SAR wind fields, sea-ice masks, as well as integrating procedures. We compared the resulting spreads in momentum flux to global values and values in the tropics, an area of prevailing low winds. We show that the choice of drag coefficient parameterization can lead to significant differences in resultant momentum flux (or wind stress) values. We found that the spread of results stemming from the choice of drag coefficient parameterization was 14 % in the Arctic, the North Atlantic and globally, but it was higher (19 %) in the tropics. On monthly time scales, the differences were larger at up to 29 % in the North Atlantic and 36 % in the European Arctic (in months of low winds) and even 50 % locally (the area west of Spitsbergen). When we choose the parameterizations which increased linearly with wind speed (7-9) momentum flux were largest for all months, in compare to values from the two parameterizations which increase with wind speed sinusoidal (12 and 13), in both regions with high and low winds and $C_D$ values were consistently higher for all wind speeds. As the one of power law parameterization (13) behaves so distinctly differently with low winds, we showed seasonal results for the tropical ocean, which were subdued for the whole region, with monthly averages in the range of 0.2 to 0.3 N m$^2$.

1. Introduction

Confirming and explaining the nature and consequence of the interaction between the atmosphere and ocean is one of the great challenge in climate and sea research. These two sphere are coupled which lead to variations covering time scales from minutes to even millennia. The purpose of this article is to examine, using a modern set of software processing tools called the FluxEngine, the nature of the fluxes of momentum across the sea surface over the North Atlantic and the European Arctic. These fluxes are important to determine of current system and sea state conditions. Our goal is to evaluate how much the average monthly and annually momentum transfer values depend on the choice of $C_D$, using the actual wind field from the North Atlantic and the European Arctic, and demonstrate existing differences as a result of the formula used.
The ocean surface mixed layer is a region where kinematic forcing affects the exchange of horizontal momentum and controls transport from the surface to depths (Gerbi et al., 2008, Bigdeli et al., 2017). Any attempt to properly model the momentum flux from one fluid to another as the drag force per unit area at the sea surface (surface shear stress, $\tau$) take into account other physical processes responsible for generating turbulence such as boundary stress, boundary buoyancy flux, and wave breaking (Rieder et al., 1994, Jones and Toba, 2001). Fluxes across the sea surface usually depend nonlinearly on the relevant atmospheric or oceanic parameters. Over the past fifty years, as the collection of flux data has increased, many empirical formulas have been developed to express the ocean surface momentum flux as a relationship between non-dimensional drag coefficient ($C_D$), wind speed ($U_{10}$), and surface roughness ($z_0$) (Wu 1969, 1982; Bunker, 1976; Garratt, 1977; Large and Pond, 1981; Trenberth et al., 1989; Yelland and Taylor, 1996, Donelan et al., 1997; Kukulka et al., 2007; Andreas et al., 2012). These formulas can be divided into two groups. One group of theories gives the $C_D$ at level $z$ in terms of wind speed and possibly one or more sea-state parameters (for example, Geernaert et al., 1987, Yelland and Taylor, 1996, Enriquez and Friehe, 1997), while the second group provides formulas for roughness length $z_0$ in terms of atmospheric and sea-state parameters (for example, Wu, 1969, Donelan et al., 1997, Andreas et al., 2012 (further referred to as A12)).

As the exchange of air-sea momentum is difficult to measure directly over the ocean meteorologist and oceanographers often rely on bulk formulas parameterized by Taylor (1916), that relate the fluxes to averaged wind speed through transfer coefficients:

$$\tau = \rho C_{Dz} U_z^2$$  \hfill (1)

where $\tau$ is the momentum flux of surface stress, $\rho$ is air density, $C_{Dz}$ is the non-dimensional drag coefficient appropriate for $z$ height, and $U_z$ is the average wind speed at some reference height $z$ above the sea. $C_{Dz}$ is commonly parameterized as a function of mean wind speed ($\text{m s}^{-1}$) for neutral-stability at a 10 m reference height above mean sea level (Jones and Toba, 2001), which is identified as $C_{DN10}$ or $C_{D10}$ (this permits avoiding deviation for the vertical flow from the logarithmic law):

$$C_{DN10} = \frac{\tau}{\rho U_{10}^2} = \left( \frac{u^*}{U_{10}} \right)^2$$  \hfill (2)

where $u^*$ is friction velocity. Alternatively, the neutrally stratified momentum flux can be determined from the logarithmic profile, thus Eq. 1 can be express as:

$$C_{DN10} = \left[ \kappa \ln \left( \frac{10}{z_0} \right) \right]^2$$  \hfill (3)

where $z_0$ (m) is the aerodynamic roughness length, which is the height, above the surface to define the measure of drag at which wind speed extrapolates to 0 on the logarithmic wind profile (Andreas et al., 2012), and $\kappa$ is von Kármán constant ($\kappa=0.4$).

At the same time, we can define the value of friction velocity by the following equation:

$$\tau = \rho u_*^2$$  \hfill (4)

Comparison with bulk formula (1) leads to the equation:

$$u_*^2 = C_{D10} U_{10}^2$$  \hfill (5)
Some of the first studies (Wu, 1969, 1982, Garrat, 1977) focused on the relationship between wind stress and sea surface roughness, as proposed by Charnock (1955), and they formulated (for winds below 15 m s\(^{-1}\)) the logarithmic dependence of the stress coefficient on wind velocity (measured at a certain height) and the von Kármán constant. Currently common parameterizations of the drag coefficient are a linear function of 10 m wind speed (\(U_{10}\)), and the parameters in the equation are determined empirically by fitting observational data to a curve. The general form is expressed as (Guan and Xie, 2004):

\[
C_D 10^3 = (a + b U_{10})
\]  

(6)

In this work our focus is on the fluxes of average values using seven different drag coefficient parameterizations (\(C_D\)), chosen for their importance for the history of the field out of many published within the last half century (Bryant and Akbar, 2016).

\[
10^3 \cdot C_{D10} = 0.5 U_{10}^{0.5}
\]  

for \(1 \text{ m s}^{-1} < U_{10} < 15 \text{ m s}^{-1}\)  

(Wu, 1969)

\[
10^3 \cdot C_{DN10} = 0.75 + 0.067 U_{10}
\]  

for \(4 \text{ m s}^{-1} < U < 21 \text{ m s}^{-1}\)  

(Garrat, 1977)

\[
10^3 \cdot C_{D10} = (0.8 + 0.065 U_{10})
\]  

for \(U_{10} > 1 \text{ m s}^{-1}\)  

(Wu, 1982)

\[
10^3 \cdot C_{DN10} = 0.29 + \frac{3.1}{U_{10N}} + \frac{7.7}{U_{10N}^2}
\]  

for \(3 \text{ m s}^{-1} < U_{10N} < 6 \text{ m s}^{-1}\)  

(Yelland and Taylor, 1996)

\[
10^3 \cdot C_{D10} = 0.60 + 0.070 U_{10N}
\]  

for \(6 \text{ m s}^{-1} < U_{10N} < 26 \text{ m s}^{-1}\)  

\[
C_{DN10} = \left(\frac{U^*_{10N}}{U_{10N}}\right)^2 = a^2 \left(1 + \frac{b}{a} U_{10N}\right)^2
\]

where \(a = 0.0583, b = -0.243\)  

(Andreas et al., 2012)

(13)

where \(C_{DN10}\) is the expression of neutral-stability (10-m drag coefficient), \(C_{D10}\) is the drag coefficient dependent on surface roughness, \(U_{10}\) is the mean wind speed measured at 10 m above the mean sea surface, \(U_{10N}\) is the 10-m, neutral-stability wind speed. All of them are generated from the vertical wind profile, but they differ in the formulas used. Two of the parameterization which we chosen are formulated as power-law of the relationship between \(C_D\) and \(U_{10}\) (eq. 7 and 13), three are formulated as linear-law (eq. 8, 9 for light winds, and 12), and one as constant value of the relationship (eq. 11). All the above studies propose different parameterizations (see Fig. 1) of the drag coefficient and the function of wind speed, which
reflects the difficulties in simultaneously measuring at high sea stress (or friction velocity) and wind speed.

Wu (1969), based on data compiled from 12 laboratory studies and 30 oceanic observations, formulated power-law (for breezes and light winds) and linear-law (for strong winds) relationships between the wind-stress coefficient ($C_y$) and wind velocity ($U_{10}$) at a certain height $y$ at various sea states. In his study, he used roughness Reynolds numbers to characterize the boundary layer flow conditions, and he assumed that the sea surface is aerodynamically smooth in the range of $U_{10} < 3 \text{ m s}^{-1}$, transient at wind speed $3 \text{ m s}^{-1} < U_{10} < 7 \text{ m s}^{-1}$, and aerodynamically rough at strong winds $U_{10} > 7 \text{ m s}^{-1}$. He also showed that the wind-stress coefficient and surface roughness increase with wind speed at light winds ($U_{10} < 15 \text{ m s}^{-1}$) and is constant at high winds ($U_{10} > 15 \text{ m s}^{-1}$) with aerodynamically rough flow. Garratt (1977), who assessed the 10 m neutral drag coefficient ($C_{DN10}$) based on 17 publications, confirmed the previous relationship and simultaneously suggested a linear form of this relationship for light wind. Wu (1980) proposed the linear-law formula for all wind velocities and later (Wu, 1982) extended this even to hurricane wind speeds. Yelland and Taylor (1996) presented results obtained from three cruises using the inertial dissipation method in the Southern Ocean and indicate that using the linear-law relationship between the drag coefficient and wind speed (for $U_{10} > 6 \text{ m s}^{-1}$) is better than using $u^*$ with $U_{10}$. The NCEP/NCAR reanalysis (Kalnay et al., 1996) uses a constant drag coefficient of $1.3 \times 10^{-3}$ while, for example, the Community Climate System Model version 3 (Collins et al., 2006) uses a single mathematical formula proposed by Large and Yeager (2004) for all wind speeds. Andreas et al. (2012) based on available datasets, friction velocity coefficient versus neutral-stability wind speed at 10 m, and sea surface roughness tested the approach proposed by Foreman and Emeis (2010) for friction velocity in order to find the best fit for parameters $a = 0.0583$ and $b = -0.243$. They justify their choice by demonstrating that $u^*$ vs. $U_{10}$ has smaller experimental uncertainty than $C_{DN10}$, and that one expression of $C_{DN10}$ for all wind speeds overstates and overestimates results in low and high winds (Figs. 7 and 8 in A12).

In this paper we investigate how the relevant or most commonly used parameterizations for drag coefficient ($C_D$) affect to value of momentum transfer values, especially in the North Atlantic (NA) and the European Arctic (EA). Our task was to demonstrate how existing differences as a result of the formula used how big they can be. As is widely known, the exact equation that describes the connection between the drag coefficient and wind speed depends on the author (Geernaert, 1990). Our intention here is not to re-invent or formulate a new drag parameterization for the NA or the EA, but to revisit the existing definition of drag parameterization, and, using satellite data, to investigate how existing formulas represent the environment in the North. We concentrated on wind speed parameterizations, because wind speed is a parameter that is available in every atmospheric circulation model. Therefore, it is used in all air-sea flux parameterizations, and presently it is used even when sea state provides a closer physical coupling to the drag coefficient (for review see Geernaert et al., 1986).

2. Materials and Methods
We calculated monthly and annual mean momentum fluxes using a set of software processing tools called the FluxEngine (Shutler et al., 2016), which was created as part of the OceanFlux Greenhouse Gases project funded by the European Space Agency (ESA). Since the toolbox, for now, is designed to calculate only air-sea gas fluxes but it does contain the necessary datasets for other fluxes, we made minor changes in the source code by adding parameterizations for the air-sea drag relationship. For the calculations, we used Earth Observation (EO) $U_{10}$ for 1992-2010 from the GlobWave project (http://globwave.ifremer.fr/). GlobWave produced a 20-year time series of global coverage multi-sensor cross-calibrated wave and wind data, which are publicly available at the Ifremer/CERSAT cloud. Satellite scatterometer derived wind fields are at present believed to be at least equally as good as wind products from reanalyses (see, for example, Dukhovskoy et al. 2017) for the area of our interest in the present study. The scatterometer derived wind values are calibrated to the $U_{10N}$, and, therefore, are fit for use with the neutral-stability drag coefficient (Chelton and Freilich, 2005). All data came in netCDF-4 format. The output data is a compilation file that contains data layers, and process indicator layers. The data layers within each output file, which are details part of the FluxEngine, include statistics of the input datasets (e.g., variance of wind speed, percentage of ice cover), while the process indicator layers include fixed masks as land, open ocean, coastal classification, and ice.

All analyses using the global data contained in the FluxEngine software produced a gridded ($1^\circ \times 1^\circ$) product. The NA was defined as all sea areas in the Atlantic sector north of $30^\circ$ N, and the EA subset was those sea areas north of $64^\circ$ N (Fig. 2). We also defined the subset of the EA east of Svalbard (“West Svalbard” between $76^\circ$ and $80^\circ$ N and 10$^\circ$ to 16$^\circ$ E), because it is a region that is studied intensively by multiple, annual oceanographic ship deployments (including that of the R/V Oceania, the ship of the institution the authors are affiliated with). FluxEngine treats areas with sea-ice presence in a way that is compatible with Lüpkes et al. (2012) multiplying the water drag coefficient by the ice-free fraction of each grid element. We also define “tropical ocean” as all areas within the Tropics ($23^\circ$ S to $23^\circ$ N, not shown) in order to test the hypothesis that the new A12 parameterization will produce significantly lower wind stress values in the region.

3. Results and Discussion

Using the FluxEngine software, we produced global gridded monthly air-sea momentum fluxes and from these we have extracted the values for the study region, the global ocean, the NA Ocean, and its subsets: the Arctic sector of the NA and the West Spitsbergen area (WS). Some of the parameterizations used were limited to a restricted wind speed domain. We used them for all the global wind speed data to avoid data gaps for winds that were too high or too low for a given parameterization (Fig. 1). However, circulation models have the very same constraint and, therefore, the procedure we used emulated using the parameterization in oceanographic and climate modeling.

Since wind velocity was used to estimate $C_D$, Fig. 1 shows a wide range of empirical formulas and Fig. 2 shows annual mean wind speed $U_{10}$ (m s$^{-1}$) in the NA and the EA. The differences between the parameterizations are distinct (Fig. 1). The $C_D$ values from the
parameterizations 7 - 9 increased linearly with wind speed since the results from the parameterizations 10,12,13 are characterized by sinusoidal distribution and indicating decreases for winds in the range of 0 - 10 m s\(^{-1}\), after which they began increase. Despite many measurements, the drag coefficient still has wide variability at low and moderate wind speeds. Our research has showed that al lower wind values (<10 m s\(^{-1}\)) the differences between the drag coefficient parameterizations are greater than at higher speeds (> 10 m s\(^{-1}\)) and the most outlier results are those obtained from the power law parameterization of Andreas et al., (2012). The lower the wind speed, the higher uncertainty are, and at low winds it is uncertainty by a factor of 0.5-1.5 depending on the formula used, while at moderate winds it is uncertainty by a factor of 1.5-2.0 (Fig. 1). At a wind value of about 15 m s\(^{-1}\), the results from eq. 8, 9, and 13 overlapped providing the same values for the drag coefficient parameterizations. Additionally, we compared directly the results of the two parameterizations for the drag air-sea relation that uses different dependencies (Fig. 1). For this estimation we chose the two most-recent parameterizations (eq. 12 and 13) that showed the lowest values and change seasonally depending on the area used. As a result, these months with weak winds have significantly lower momentum flux values, which could be the effect of statistically weaker wind in ocean areas having stable winds with waves traveling in the same direction as the wind at similar speeds.

Comparison showed that the A12 parameterization demonstrates almost zero sea surface drag for winds in the range of 3 - 5 m s\(^{-1}\), which is compensated for by a certain surplus value for strong winds. The small drag coefficient values facilitate what Grachev and Fairall (2001) describe as the transfer of momentum from the ocean to the atmosphere at wind speeds of 2 - 4 m s\(^{-1}\), which correspond to the negative drag coefficient value. Such events require specific meteorologist conditions, but this strongly suggests that the average \(C_D\) value for similar wind speeds could be close to zero. The annual mean wind speed in the NI is 10 m s\(^{-1}\), and in the EA it is 8.5 m s\(^{-1}\) (Fig. 2).

Figure 3 presents maps of the mean boreal winter DJF and summer JJA momentum fluxes for the chosen \(C_D\) parameterizations (Wu, 1969 and A12 – the ones with the largest and smallest \(C_D\) values). The supplementary materials contain complete maps of annual and seasonal means for all the parameterizations. The zones of the strongest winds are in the extratropics in the winter hemisphere (southern for JJA and northern for DJF). The older Wu (1969) parameterization produces higher wind stress values than A12 in both regions with high and low winds and \(C_D\) values are consistently higher for all wind speeds except the lowest ones (which, after multiplying by \(U^2\), produced negligible differences in wind stress for the lowest winds). The average monthly values for each of the studied areas are shown in Fig. 4. Generally, this illustrates that the sinusoidal the drag coefficient parameterization is, the smaller the calculated momentum flux is. For global data (Fig. 4a), not much seasonal change is noted, because the strongest winds are in fall and winter, but these seasons are the opposite in the northern and southern hemispheres. The parameterization with the largest momentum flux values for all months is that of Wu (1969), the linear one, while the two parameterizations with the lowest values are the sinusoidal ones (Large and Yeager, 2004 and A12). For the NI (Fig. 4b), with is much more pronounced seasonal wind changes, the situation is more complicated. With high winter winds, the A12 parameterization is no longer the one that produces the smallest wind stress (it is actually in the middle of the seven). However, for low summer winds,
it is the lowermost outlier. Actually, in summer, the constant $C_p$ value used by the NCEP/NCAR reanalysis produces the highest wind stress values in the NA. The situation is similar for the EA (a subset of the NA), the wind stress values of which are shown in Fig. 4c, and for the WS area (not show). In the Arctic summer, A12 produces the least wind stresses, while all the other parameterizations look very similar qualitatively (even more so in the Arctic than in the whole NA). Because the A12 parameterization behaves so distinctly differently with low winds, we also show seasonal results for the tropical ocean (Fig. 4d). The seasonal changes are subdued for the whole tropical ocean with the slight domination of the Southern Hemisphere (the strongest winds are during the boreal summer) with generally lower momentum transfer values (monthly averages in the range of 0.2 to 0.3 N m$^{-2}$ compared to 0.2 to 0.4 N m$^{-2}$ for the NA and 0.2 to 0.5 N m$^{-2}$ for the Arctic). The sequence of values for the parameterization is similar to that of the global ocean, but there are differences. Here the NCEP/NCAR constant parameterization is the second highest (instead of Wu, 1982 for the global ocean) while, unlike in the case of the global ocean, A12 produces visibly lower values than does the Large and Yeager (2004) parameterization.

Table 1 and Fig. 5 present the annual average air-sea momentum flux values (in N m$^{-2}$) for all the all regions studied and all the parameterizations. The results show that the annual North Atlantic momentum fluxes, depending on the formula used, varies from -0.0.290 N m$^{-2}$ for A12 to 0.333 N m$^{-2}$ for Wu (1969). In the case of global annual average, the values are $0.283$ and $0.322$, respectively. Table 1 shows also the same data “normalized” to the A12 data (presented as percentages of A12, which produced the lowest values for each region), which allows us to visualize the relative differences. A surprising result is the annual ratios of the parameterizations values for the global, the NA, and the Arctic regions (Fig. 4 shows that this is not true on monthly scales). The spread of the momentum flux results is 14 % in all three regions, and even flux values themselves are larger in the NA than globally and larger in the Arctic than in the whole of the NA basin. In the NA region with winds stronger than average for world ocean, the formula giving highest momentum transfer results are the ones with highest values for strong winds, with exception of Andreas et al. (2012) which is lower due to its low values for lower winds speeds. The smaller WS region, with winds that are, on average, weaker than those of the whole Arctic (but stronger than those of the whole NA), had slightly different ratios of the resultant fluxes. For the tropical ocean, which is included for comparison because of its weaker winds, the spread in momentum flux values on an annual scale is 19 %. The spreads are even larger on monthly scales (not shown). The difference between A12 and Wu (1969) and NCEP/NCAR (the two parameterizations producing the largest fluxes on monthly scales) are 27 % and 29 % for the NA (in July), 31 % and 36 % for the Arctic (in June), 42 % and 51 % for the WS region (in July) and 23 % and 22 % for the tropical ocean (in April), respectively. Seasonality in the tropics is weak, therefore, the smallest monthly difference of 16 % (July) is larger than the difference for the global data in any month (the global differences between the parameterizations have practically no seasonality). On the other hand, the smallest monthly differences between the parameterizations in the NA, the Arctic, and the WS regions are all 7 %, in the month of the strongest winds (January).
Because the value of momentum flux is important for ocean circulation, its correct calculation in coupled models is very important, especially in the Arctic, where cold halocline stratification depends on the amount of mixing (Fer, 2009). We show that with the parameterization used in modelling, such as the NCEP/NCAR constant parameterization and Large and Yeager (2004), production stress results differ by about 5 %, on average (both in the Arctic and globally), and the whole range of parameterizations leads to results that differ, on average, by 14 % (more in low wind areas) and much more on monthly scales. One aspect that needs more research is the fact that the newest parameterization, A12, produces less momentum flux than all the previous ones, especially in lower winds (which, by the way, continues the trend of decreasing values throughout the history of the formulas discussed). The A12 parameterization is based on the largest set of measurements of friction velocity as a function of wind speed and utilizes the recently discovered fact that $b$ in equation (7) is not negligible. It also fits the observations that developed swell at low wind velocity has celerity which leads to zero or even negative momentum transfer (Grachev and Fairall, 2001). Therefore, the significantly lower A12 results for the tropical ocean (the trade wind region) and months of low winds elsewhere could mean that most momentum transfer calculations are overestimated. This matter needs further study, preferably with new empirical datasets.

4. Conclusions

In the present work the evaluation of how the selected parameterization affects the total value of momentum fluxes for large reservoirs was assessed. This allows constraining the uncertainty caused by the parameterization choice. In order to achieve this we calculated monthly and annually average momentum fluxes using a set of software processing tools called the FluxEngine in the North Atlantic (NA) and the European Artic (EA). The NA was defined as all sea surface areas in the Atlantic sector north of 30° N, and the EA was sea areas north of 64° N. Based on our results, we still do not know which one of the parameterizations can be recommended as the most suitable for the NA and the EU study. Further investigation of the differences in the parameterization of the exchange coefficient in the various algorithm would help in resolving this problem.

Despite many measurements, the drag coefficient still has wide variability at low and moderate wind speeds. The lower the wind speed, the higher uncertainty are, and at low winds it is uncertainty by a factor of 0.5-1.5 depending on the formula used, while at moderate winds it is uncertainty by a factor of 1.5-2.0 (Fig. 1). The annual mean wind speed in the NA is 10 m s$^{-1}$, and in the EU it is 8.5 m s$^{-1}$.

We show that the choice of drag coefficient parameterization can lead to significant differences in resultant momentum flux (or wind stress) values. Comparing the values of momentum flux across the sea surface from the power law parameterization, it showed that in both regions, with low and high winds, the parameterizations specified for all winds speeds (eq. 13) has lower values of wind stress than the parameterizations specified for light winds (eq. 7). In the Arctic, the NA, and globally the differences between the wind stress, depend on formula used, are 14 % and they are higher in low winds areas. The parameterizations generally have a decreasing trend in the resultant momentum flux values, with the most recent (Andreas et al., 2012) producing the lowest wind stress values, especially at low winds, resulting in almost 20
% differences in the tropics (Table1). The differences can be much larger on monthly scales, up to 29 % in the NA and 36 % in the EA (in months of low winds) and even 50 % locally in the area west of Spitsbergen. For months that have the highest average winds, the percentage differences are smaller (about 7 % everywhere), but because absolute value of the flux are largest for high winds, this 7% discrepancy is also important for air-sea momentum flux values.

Since momentum flux is an important parameter in ocean circulation modeling, we believe more research is needed (one aspects that needs more research is the fact that the newest power law parameterization, A12, produces less momentum flux than all the previous ones, especially in lower winds), and the parameterizations used in the models possibly need further development.

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References


Table 1. Area average annual mean values of momentum flux (wind stress) [N m⁻²] for all the studied regions and parameterizations. In each column the percentage values are normalized to A12, the parameterization that produced the smallest average flux values.
Figure 1. The drag coefficient parameterization used in the study (Eqs. 7-13) as a function of wind speed $U_{10}$ (m s$^{-1}$).

Figure 2. Annual mean wind speed $U_{10}$ (m s$^{-1}$) in the study area—the North Atlantic and the European Arctic (north of the red line).

Figure 3. Maps of momentum flux [N m$^{-2}$] across the sea surface (wind stress) for boreal winters ((a) and (c)) and summers ((b) and (d)) for Wu (1969) and A12 drag coefficient parameterizations (the two parameterizations with the highest and lowest average values, respectively).

Figure 4. Monthly average momentum flux values [N m$^{-2}$] for (a) global ocean, (b) North Atlantic, (c) European Arctic, and (d) tropical ocean. The regions are defined in the text.

Figure 5. Area annual average momentum flux values for (a) European Arctic and (b) Tropical ocean. The vertical solid line is the average of all seven parameterization and the dashed lines are standard deviations for the presented values. Global and the North Atlantic results are not shown because the relative values for different parameterizations are very similar (see Table 1), scaling almost identically between the basins.
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<table>
<thead>
<tr>
<th></th>
<th>Global</th>
<th>North Atlantic</th>
<th>Arctic</th>
<th>W. Spitsbergen</th>
<th>Tropics</th>
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<tbody>
<tr>
<td>Wu (1969)</td>
<td>0.322 (114%)</td>
<td>0.330 (114%)</td>
<td>0.375 (114%)</td>
<td>0.360 (114%)</td>
<td>0.261 (119%)</td>
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<td>Garratt (1977)</td>
<td>0.307 (109%)</td>
<td>0.316 (109%)</td>
<td>0.358 (109%)</td>
<td>0.344 (110%)</td>
<td>0.251 (115%)</td>
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<tr>
<td>Wu (1982)</td>
<td>0.311 (110%)</td>
<td>0.320 (110%)</td>
<td>0.363 (110%)</td>
<td>0.349 (111%)</td>
<td>0.255 (117%)</td>
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<tr>
<td>NCEP/NCAR</td>
<td>0.303 (107%)</td>
<td>0.312 (107%)</td>
<td>0.353 (107%)</td>
<td>0.341 (108%)</td>
<td>0.258 (118%)</td>
</tr>
<tr>
<td>Yelland &amp; Taylor (1996)</td>
<td>0.297 (105%)</td>
<td>0.306 (105%)</td>
<td>0.348 (106%)</td>
<td>0.335 (107%)</td>
<td>0.245 (112%)</td>
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<tr>
<td>Large &amp; Yeager (2004)</td>
<td>0.285 (101%)</td>
<td>0.293 (101%)</td>
<td>0.333 (101%)</td>
<td>0.320 (102%)</td>
<td>0.236 (108%)</td>
</tr>
<tr>
<td>Andreas et al., (2012)</td>
<td>0.283 (100%)</td>
<td>0.290 (100%)</td>
<td>0.329 (100%)</td>
<td>0.314 (100%)</td>
<td>0.219 (100%)</td>
</tr>
</tbody>
</table>
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(a) Wu (1969)  
(b) Wu (1969)


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North Atlantic mean momentum flux
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